

Tutorials for the course “Relativitätstheorie und Kosmologie I”: Problem Sheet 10

Please follow the instructions on moodle concerning which problems should be used for the Kreuzerlliste.

We write $\det g_{\mu\nu}$ for the determinant of a matrix $(g_{\mu\nu})$, $g^{\mu\nu}$ for the matrix inverse to $g_{\alpha\beta}$ (thus, $g^{\mu\nu}g_{\nu\alpha} = \delta^\mu_\alpha$), and $\phi^A_{,\mu}$ for $\partial\phi^A/\partial x^\mu$.

47 Let $A_{\mu\nu}$ be an invertible matrix, with inverse $B^{\mu\nu}$. Show that

$$\frac{\partial(\det A_{\mu\nu})}{\partial A_{\rho\sigma}} = (\det A_{\mu\nu})B^{\rho\sigma}. \quad (1)$$

Deduce that for matrices with entries which are functions on space-time

$$\partial_\alpha \sqrt{\det A_{\mu\nu}} = \frac{1}{2} \sqrt{\det A_{\mu\nu}} B^{\rho\sigma} \partial_\alpha A_{\rho\sigma}. \quad (2)$$

48 Write-down the Euler-Lagrange equations, the canonical energy-momentum tensor

$$t^\mu{}_\nu = \frac{\partial \mathcal{L}}{\partial \phi^A_{,\mu}} \phi^A_{,\nu} - \mathcal{L} \delta^\mu_\nu,$$

and the symmetric energy-momentum tensor

$$T_{\mu\nu} = 2 \frac{\partial \mathcal{L}}{\partial g^{\mu\nu}},$$

for the following Lagrangeans:

- i. $\mathcal{L} = \frac{1}{2} (g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi + m^2 \phi^2) \sqrt{-\det g_{\mu\nu}}$
(massive scalar field);
- ii. $\mathcal{L} = \left(\frac{1}{8\pi} g^{\alpha\beta} g^{\mu\nu} F_{\alpha\mu} F_{\beta\nu} \right) \sqrt{-\det g_{\mu\nu}}$,
where $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ (Maxwell field); and
- iii. $\mathcal{L} = \frac{1}{8\pi} (e^{-a\phi} g^{\alpha\beta} g^{\mu\nu} F_{\alpha\mu} F_{\beta\nu} + g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi) \sqrt{-\det g_{\mu\nu}}$,
where as before $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ and $a \in \mathbb{R}$ (dilaton-Maxwell system).

In the calculations of the Euler-Lagrange equations and of the canonical energy-momentum tensor you can assume that $g_{\mu\nu} = \eta_{\mu\nu}$, the Minkowski metric.

49 Changes of coordinates

Let $x^0 = t$, $x^1 = x$, $x^2 = y$, $x^3 = z$ be inertial coordinates on flat space-time, so the Minkowski metric has components

$$(g_{ab}) = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

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Let X be the vector field which in the above coordinate system equals $(1, 1, 0, 0)$, and let α be a one-form which in the above coordinate system equals $(1, 1, 0, 0)$.

Find the metric coefficients \tilde{g}_{ab} , and the components of X and α , in each of the following coordinate systems.

(i) $\tilde{x}^0 = t - z, \tilde{x}^1 = r, \tilde{x}^2 = \theta, \tilde{x}^3 = z$

(ii) $\tilde{x}^0 = \tau, \tilde{x}^1 = \phi, \tilde{x}^2 = y, \tilde{x}^3 = z,$

(iii) $\tilde{x}^0 = t, \tilde{x}^1 = r, \tilde{x}^2 = \theta, \tilde{x}^3 = \phi$

where, in the first case, r, θ are plane polar coordinates in the x, y plane: $x = r \cos \theta$, $y = r \sin \theta$; in the second, τ, ϕ are ‘Rindler coordinates’, defined by $t = \tau \cosh \phi$, $x = \tau \sinh \phi$; and, in the third, r, θ, ϕ are spherical polar coordinates. In each case, state which region of Minkowski space the coordinate system covers. *[Hint: A quick method for the metric is to write it as $ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$ and substitute, for example, $dx = \cos \theta dr - r \sin \theta d\theta$, and so on. Of course you should convince yourself that this is legitimate.]*

50 Lie bracket

Recall that vector fields can be identified with homogeneous linear first order partial differential operators $X = X^a \partial_a$ acting on functions as $X(f) = X^a \partial_a f$.

The Lie-bracket $[X, Y]$ of two vector fields X and Y is defined as

$$[X, Y](f) = X(Y(f)) - Y(X(f)) .$$

Show that $[X, Y]$ also is a vector field, i.e. a homogeneous linear first order differential operator, with components

$$[X, Y]^a = X^b \partial_b Y^a - Y^b \partial_b X^a . \tag{3}$$

Check, by a direct coordinate calculation, that the right-hand-side of (3) transforms as a vector field under changes of coordinates.

[For self-study:] Prove the *Jacobi identity*:

$$[X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0 .$$