

Übungen zur Vorlesung Relativitätstheorie und Kosmologie II: Problem Sheet 4

- 1 Some timelike geodesics in Schwarzschild. Let $V^2 = 1 - 2m/r$, and consider the Schwarzschild metric:

$$g = -\left(1 - \frac{2m}{r}\right)dt^2 + \frac{dr^2}{1 - \frac{2m}{r}} + r^2 d\Omega^2, \quad (1)$$

$$t \in \mathbb{R}, r \neq 2m, 0, \quad (2)$$

Recall, from a previous problem sheet, that for geodesics in the Schwarzschild metric lying in the equatorial plane $\theta = \pi/2$ we have the following constants of motion (you are strongly encouraged to check that you can reproduce the result, but this will not be done again in class):

$$\frac{d}{ds} \left(V^2 \frac{dt}{ds} \right) = 0 \implies \frac{dt}{ds} = \frac{E}{1 - \frac{2m}{r}}. \quad (3)$$

$$\frac{d}{ds} \left(r^2 \frac{d\varphi}{ds} \right) = 0 \implies \frac{d\varphi}{ds} = \frac{J}{r^2}. \quad (4)$$

$$\underbrace{V^2 \left(\frac{dt}{ds} \right)^2}_{E^2 V^{-2}} - V^{-2} \left(\frac{dr}{ds} \right)^2 - \underbrace{r^2 \left(\frac{d\varphi}{ds} \right)^2}_{J^2/r^2} = \lambda \in \{0, \pm 1\}. \quad (5)$$

Consider a timelike geodesic in the Schwarzschild metric parameterized by proper time, thus $\lambda = 1$, and lying in the equatorial plane $\theta = \pi/2$.

Q1. Verify that

$$\frac{E^2 - \dot{r}^2}{1 - \frac{2m}{r}} - \frac{J^2}{r^2} = 1.$$

Q2. Deduce that if $E = 1$ and $J = 4m$ then

$$\frac{\sqrt{r} - 2\sqrt{m}}{\sqrt{r} + 2\sqrt{m}} = A e^{\epsilon\varphi/\sqrt{2}},$$

where $\epsilon = \pm 1$ and A is a constant. Describe the orbit that starts at $\varphi = 0$ in each of the cases (i) $A = 0$, (ii) $A = 1$, $\epsilon = -1$, (iii) $r(0) = 3m$, $\epsilon = -1$.

Q3: Consider the case $E = 1$, $J = 0$ and deduce a relation for $r = r(s)$.

- 2 Null Geodesics in Schwarzschild Let γ be an affinely parameterized null geodesic, thus $\lambda = 0$, in the Schwarzschild metric lying in the equatorial plane $\theta = \pi/2$. Assume that $J \neq 0$, hence we can make a change of parameter $\varphi \mapsto s(\varphi)$ using the implicit equation

$$\frac{d\varphi}{ds} = \frac{J}{r^2} \implies \frac{ds}{d\varphi} = \frac{r^2}{J}.$$

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Set

$$u(\varphi) = \frac{m}{r(s(\varphi))}, \quad p(\varphi) = \frac{du(\varphi)}{d\varphi}.$$

Q1. Show that along γ , for $u \neq \frac{1}{2}$, we have

$$p^2 = 2u^3 - u^2 + \alpha^2. \quad (6)$$

Q2. For the case $\alpha = 0$ show that we obtain an autonomous two-dimensional system

$$\frac{du}{d\varphi} = p, \quad \frac{dp}{d\varphi} = 3u^2 - u.$$

Recall that critical points of a dynamical system $d\vec{x}/d\varphi = \vec{Y}$ (here $\vec{x} = (u, p)$) are defined as points where \vec{Y} vanishes. Find the critical points. Can you sketch the trajectories in the (u, p) phase-plane?

Q3. For the case $\alpha = 0$ and $u > 1/2$ show that the geodesic has an equation of the form

$$r = 2m \cos^2\left(\frac{\varphi - \varphi_0}{2}\right), \quad (7)$$

with $t = t(\varphi)$ that you should determine.