

- 1 Let a bracket over indices denote complete antisymmetrisation, and a parenthesis over indices denote complete symmetrisation: for example,

$$A_{[\mu\nu]} = \frac{1}{2}(A_{\mu\nu} - A_{\nu\mu}), \quad A_{(\mu\nu)} = \frac{1}{2}(A_{\mu\nu} + A_{\nu\mu}), \quad \delta_{\mu}^{[\alpha} \delta_{\rho}^{\gamma]} = \frac{1}{2}(\delta_{\mu}^{\alpha} \delta_{\rho}^{\gamma} - \delta_{\mu}^{\gamma} \delta_{\rho}^{\alpha}),$$

and similarly for 3, 4 or more indices. Show that

- i. $A^{[\mu\nu]} B_{\mu\nu} = A^{\mu\nu} B_{[\mu\nu]}$,
- ii. $A^{[\mu\nu\rho]} B_{\mu\nu\rho} = A^{\mu\nu\rho} B_{[\mu\nu\rho]}$,
- iii. $\delta_{\mu}^{[\alpha} \delta_{\rho}^{\gamma]} = \delta_{[\mu}^{\alpha} \delta_{\rho]}^{\gamma}$,
- iv. $\delta_{\mu}^{[\alpha} \delta_{\nu}^{\beta} \delta_{\rho}^{\gamma]} = \delta_{[\mu}^{\alpha} \delta_{\nu}^{\beta} \delta_{\rho]}^{\gamma}$,
- v. $\epsilon^{\alpha\beta\gamma\delta} \epsilon_{\alpha\beta\gamma\rho} = -6\delta_{\rho}^{\delta}$,
- vi. $\epsilon^{\alpha\beta\gamma\delta} \epsilon_{\alpha\beta\nu\rho} = -4\delta_{[\nu}^{\gamma} \delta_{\rho]}^{\delta}$.

- 2 Assuming the tensorial transformation law of $F^{\mu\nu}$, derive the explicit formulae for the transformation laws of the electric and magnetic fields under a boost along the first coordinate axis.
- 3 1) Assuming that Λ^{α}_{β} is a Lorentz matrix, show that $\eta^{\alpha\beta} \Lambda^{\mu}_{\beta} \eta_{\mu\nu}$ is inverse to Λ^{α}_{β} .
2) Recall that we required that $F^{\mu\nu}$ transforms as follows under Lorentz transformations: if $\bar{x}^{\alpha} = \Lambda^{\alpha}_{\beta} x^{\beta} + a^{\alpha}$, then

$$\bar{F}^{\mu\nu}(\bar{x}) = \Lambda^{\mu}_{\alpha} \Lambda^{\nu}_{\beta} F^{\alpha\beta}(x),$$

and that $F_{\mu\nu}$ has been defined as

$$F_{\mu\nu} := \eta_{\mu\alpha} \eta_{\nu\beta} F^{\alpha\beta}.$$

Use 1) to show that $F_{\alpha\beta}$ transforms as

$$\bar{F}_{\mu\nu}(\bar{x}) = (\Lambda^{-1})^{\alpha}_{\mu} (\Lambda^{-1})^{\beta}_{\nu} F_{\alpha\beta}(x)$$

(this is called the *transformation law of a two-covariant tensor*);

- 4 Show that contractions $F_{\alpha\beta} F^{\alpha\beta}$ and $F^{\alpha\beta} *F_{\alpha\beta}$ are invariant (more precisely, behave as scalars) under Lorentz transformations. Express them in terms of \vec{E} and \vec{B} .
- 5 Let $\vec{E} \cdot \vec{B} = 0$, and suppose that $|\vec{E}|^2 \neq |\vec{B}|^2$. Show that there exists a Lorentz frame in which either \vec{E} or \vec{B} vanishes.